

# Comparative Study on Light-Induced Bias Stress Instability of IGZO Transistors With $\text{SiN}_x$ and $\text{SiO}_2$ Gate Dielectrics

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**Abstract**—This letter examines the effect of the gate dielectric material on the light-induced bias-temperature instability of an In–Ga–Zn–O (IGZO) thin-film transistor (TFT). After applying positive and negative bias stresses, the  $\text{SiN}_x$ -gated TFT exhibited inferior stability to the  $\text{SiO}_2$ -gated TFT, which was explained by the charge trapping mechanism. However, light illumination under a negative bias stress accelerated the negative displacement of the threshold voltage ( $V_{\text{th}}$ ) of the  $\text{SiN}_x$ -gated IGZO TFT compared to that of the  $\text{SiO}_2$ -gated TFT. This was attributed to the injection of photocreated hole carriers into the underlying gate dielectric bulk region as well as the hole trapping at the gate/channel interface.

**Index Terms**—a-InGaZnO, amorphous semiconductor, bias stability, multicomponent oxide semiconductor, thin-film transistors (TFTs).

## I. INTRODUCTION

RECENTLY, In–Ga–Zn–O (IGZO) thin-film transistors (TFTs) have attracted considerable interest for applications to the active-matrix (AM) backplane of LCD and OLED displays [1]–[4]. Although high-performance IGZO TFTs have been fabricated, light- and bias-induced instability is still a critical issue to be resolved because the exposure of the switching transistor to backlight or ambient light would be inevitable during the operation of the display panel. In this regard, there have been several reports, including the effect of moisture [5], device configuration [6], and channel cation composition [7], on the light-induced bias instability. However, there are no reports of the effect of the gate dielectric material on the photoreliability of IGZO TFTs. This letter reports the degradation mechanism

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of the threshold voltage ( $V_{\text{th}}$ ) under photon-irradiated gate bias stress conditions for IGZO TFTs. In particular, the impact of the different gate dielectrics on the light-induced bias-thermal instability of IGZO TFTs was investigated. The fabricated device had  $\text{SiN}_x$  and  $\text{SiO}_2$  films as a gate dielectric due to their large area scalability in flat panel AM displays.

## II. EXPERIMENTAL PROCEDURE

A 200-nm-thick  $\text{SiN}_x$  film as a gate dielectric layer was deposited by plasma enhanced chemical vapor deposition (PECVD) on a patterned Mo (2000-nm-thick)/glass substrate at 380 °C. An a-IGZO film, 50 nm in thickness, was grown by dc sputtering at room temperature. After defining the IGZO channel using photolithography and wet etching, a 200-nm-thick  $\text{SiO}_2$  etch stopper was deposited by PECVD and patterned by dry etching to define the source/drain contact hole. A Mo source/drain electrode (200 nm) was formed using a conventional sputtering system and patterned by photolithography and dry etching.

The fabricated TFTs had a bottom-gate and top-contact configuration [3]. For comparison, a PECVD-derived  $\text{SiO}_2$  film as a gate insulator was also deposited before active channel deposition. The thickness of the  $\text{SiO}_2$  gate insulator was designed to be 120 nm to have a similar gate capacitance: The gate capacitances per unit area of the  $\text{SiN}_x$ - and  $\text{SiO}_2$ -gated devices were  $2.66 \times 10^{-8}$  and  $2.88 \times 10^{-8} \text{ F/cm}^2$ , respectively. Finally, the samples were subject to thermal annealing at 350 °C for 1 h. The electrical measurements were carried out in air using an Agilent B1500 A precision semiconductor parameter analyzer.

## III. RESULTS AND DISCUSSION

Fig. 1(a) and (b) shows the transfer characteristics of the IGZO TFTs with  $\text{SiN}_x$  (device A) and  $\text{SiO}_2$  (device B) as a gate insulator ( $W/L = 24/10 \mu\text{m}$ ), respectively. The  $\mu_{\text{FE}}$  value was calculated from the maximum transconductance ( $\mu_{\text{FE}} = Lg_m/WC_iV_{\text{DS}}$ , where  $C_i$  and  $g_m$  are the gate capacitance per unit area and the transconductance, respectively) [4]. The threshold voltage ( $V_{\text{th}}$ ) was determined by adjusting the gate voltage, which induces a drain current of  $L/W \times 10 \text{ nA}$  at a  $V_{\text{DS}}$  of 15.1 V. In addition, the subthreshold gate swing [ $SS = dV_{\text{GS}}/d \log I_{\text{DS}}$  (V/dec)] was extracted from the linear portion of the  $\log(I_{\text{DS}})$ -versus- $V_{\text{GS}}$  plot. Device A exhibited a  $\mu_{\text{FE}}$  of

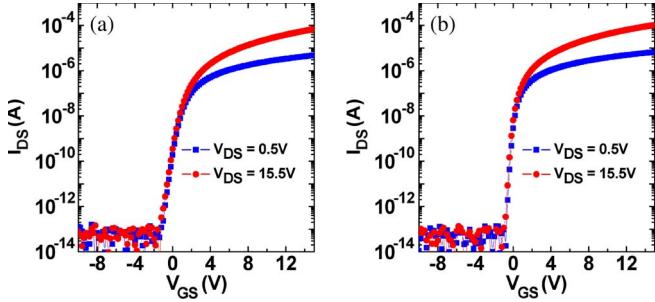


Fig. 1. Transfer characteristics of a-IGZO TFTs with a bottom-gate and top-contact structure ( $W/L = 24/10 \mu\text{m}$ ). (a)  $\text{SiN}_x$ -gated device (device A). (b)  $\text{SiO}_2$ -gated device (device B).

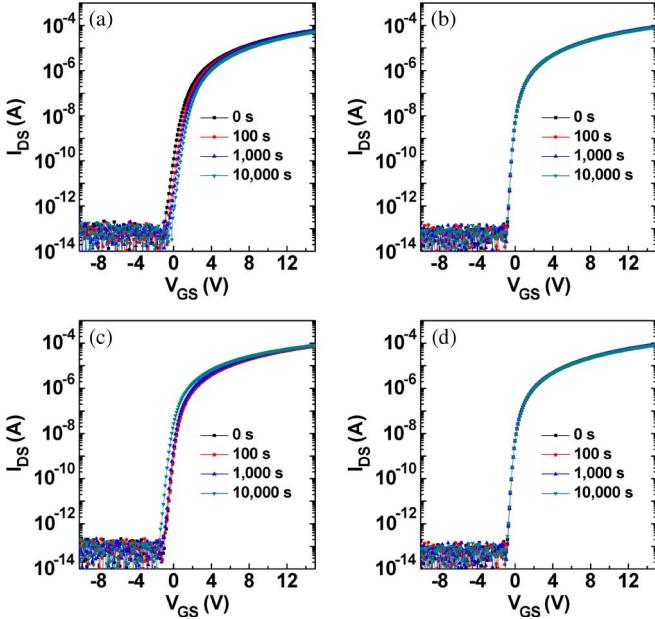


Fig. 2. Evolution of the transfer characteristics of (a) device A and (b) device B with increasing PBS time under the dark condition. Evolution of the transfer curves of (c) device A and (d) device B with increasing NBS time under the dark condition.

$13.7 \pm 0.08 \text{ cm}^2/\text{V} \cdot \text{s}$ ,  $SS$  value of  $0.38 \pm 0.01 \text{ V/dec}$ ,  $V_{th}$  of  $0.7 \pm 0.17 \text{ V}$ , and  $I_{ON}/OFF$  ratio of  $> 10^9$ . On the other hand, device B showed significantly improved performance: The  $\mu_{FE}$  and  $SS$  values were enhanced to  $16.9 \pm 0.12 \text{ cm}^2/\text{V} \cdot \text{s}$  and  $0.16 \pm 0.006 \text{ V/dec}$ , respectively, with a similar  $V_{th}$  ( $0.1 \text{ V}$ ) and  $I_{ON}/OFF$  ratio ( $> 10^9$ ). The device performance reported herein is comparable to the state-of-the-art characteristics for the IGZO TFTs with various novel gate insulators, including  $\text{Y}_2\text{O}_3$  [2],  $\text{Al}_2\text{O}_3$  [8], and  $\text{ZrO}_2$  [9].

Fig. 2(a) and (b) shows the evolution of the transfer curve as a function of the applied positive gate bias stress (PBS) time for devices A and B, respectively. The devices were stressed under the following conditions:  $V_{GS}$  was set to  $20 \text{ V}$  at room temperature, and the stress duration was  $10000 \text{ s}$ . For device A, a parallel  $V_{th}$  shift of  $0.77 \text{ V}$  to a higher voltage with increasing PBS time occurred without an accompanying change in  $\mu_{FE}$ ,  $SS$ , and the  $I_{ON}/OFF$  ratio. This positive  $V_{th}$  shift after the application of the PBS has been attributed frequently to the charge trapping of accumulated electron carriers in the channel region at/near the gate dielectric and channel interface [10], [11]. The ambient effect, such as oxygen or water molecules,

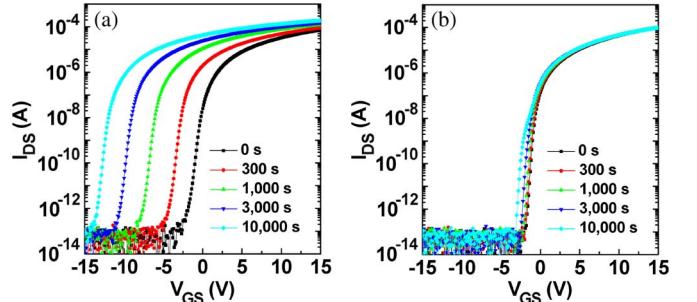


Fig. 3. Evolution of the transfer characteristics of (a) device A and (b) device B with increasing NBS time under white light illumination.

can be excluded because high-quality  $\text{SiO}_2$  films (etch stop layer) were coated on the channel region as a passivation layer [12]. The  $V_{th}$  stability of the IGZO TFTs under the PBS was improved dramatically by replacing the gate insulator with a PECVD-derived  $\text{SiO}_2$  film [see Fig. 2(b)]. A  $V_{th}$  shift of only  $6 \text{ mV}$  was observed for device B after an identical PBS application for  $10000 \text{ s}$ . If it is assumed that the  $V_{th}$  instability of the IGZO TFTs is due to the charge trapping at/near the gate/channel interface, the interfacial trap density ( $N_{it}$ ) should adversely affect the bias stability. The maximum  $N_{it}$  value can be extracted from the  $SS$  value ( $N_{it} = (SS \log(e)/(kT/q) - 1)C_i/q$ , where  $k$  is the Boltzmann's constant,  $T$  is the absolute temperature, and  $q$  is the elementary electron charge) [13], which assumes that there is no semiconductor bulk trap in the IGZO thin film. The calculated  $N_{it,max}$  ( $3.0 \times 10^{11}/\text{cm}^2$ ) for device B was much smaller than that ( $8.9 \times 10^{11}/\text{cm}^2$ ) for device A. This lower  $N_{it,max}$  for device B would explain why only a negligible  $V_{th}$  shift caused by the charge trapping was observed for device B. The effect of the gate dielectric material on the negative gate bias instability of the resulting IGZO TFTs was examined. The devices were stressed under the following conditions:  $V_{GS}$  was set to  $-20 \text{ V}$ , and the stress duration was  $10000 \text{ s}$ . The application of a negative gate bias stress (NBS) to devices A and B at room temperature caused negative  $V_{th}$  shifts of  $0.26$  and  $0.20 \text{ V}$ , respectively. At an elevated temperature of  $60^\circ\text{C}$ , the negative movements of  $V_{th}$  for devices A and B were increased to  $0.53$  and  $0.29 \text{ V}$ , respectively, as shown in Fig. 2(c) and (d). Therefore, the NBS-induced instability of device A was inferior to that of device B at  $60^\circ\text{C}$  and room temperature. This is consistent with the aforementioned charge trapping model. Hole carriers, which may be created in the channel region under NBS conditions, tend to be trapped at the interfacial defect states. Therefore, device A with a larger  $N_{it}$  would suffer from a larger charge trapping phenomenon leading to enhanced  $V_{th}$  instability.

To double check this interpretation, the TFTs were irradiated with white light during the application of the NBS ( $V_{GS} = -20 \text{ V}$ ). A white halogen lamp was used as a light source, which consisted of visible light ranging from  $400\text{-}700\text{-nm}$  wavelengths. The intensity of the light source was  $\sim 1 \text{ mW/cm}^2$ , which was calibrated by photometry. The intensity of  $1 \text{ mW/cm}^2$  was chosen to meet the required brightness for the practical application of flat panel displays [14]. Fig. 3(a) and (b) shows the time evolution of the transfer curves for both devices as a function of the light-illuminated NBS condition

for devices A and B, respectively. The inferior stability of the nitride device was manifested more clearly under the light-illuminated NBS condition. Although the  $V_{th}$  value for device A was shifted significantly to the negative direction by 11.6 V, device B exhibited only  $V_{th}$  movement of  $-0.78$  V. The consistent inferior stability of the TFT with the  $\text{SiN}_x$  gate dielectric is believed to originate from its higher  $N_{it}$  value, which is reconciled with the aforementioned charge trapping model. Under the light-illuminated NBS condition, electron-hole pairs in the channel region would be generated by light. The photocreated electron carrier will be repelled to the back channel surface, whereas the photogenerated hole carriers are attracted to the gate dielectric layer by the negative gate voltage and become trapped at the interfacial trap state. This has been suggested as a plausible explanation for the light-enhanced negative bias instability of oxide TFTs [5], [6]. However, it is not obvious that the approximately 14-fold increase in the amount of the  $V_{th}$  movement for device A (device A:  $\Delta V_{th} = -11.6$  V; device B:  $\Delta V_{th} = -0.78$  V) under the light-illuminated NBS condition can be attributed entirely to the approximately three times higher  $N_{it,\max}$  of device A. Indeed, the hole carriers near the channel region feel the hole potential barrier (i.e., valence band offset) at the interface between the gate dielectric and channel layer. Interestingly, the valence offset between the  $\text{SiN}_x$  and IGZO film can be estimated as  $\sim 0.15$  eV, which is much smaller than ( $\sim 2.80$  eV) that between the  $\text{SiO}_2$  and IGZO film [15], [16]. Therefore, the photocreated hole carriers can be injected easily into the  $\text{SiN}_x$  gate dielectric due to its lower valence band offset value, which would be partly responsible for the inferior stability of device A. Therefore, the possibility of hole injection into the gate dielectric bulk region cannot be excluded. It is noted that the fast recovery of  $V_{th}$  was observed for device A. Without any thermal annealing, device A exhibited a significant recovery of  $V_{th}$  by  $\sim 8$  V after the 10-h exposure to the air ambient, suggesting that most of the photocreated trap states ( $\sim 70\%$ ) are temporal traps. On the other hand, the  $250^\circ\text{C}$  thermal annealing was found to result in the full recovery of  $V_{th}$ , which indicates that the hole carrier injection indeed cannot be neglected.

Finally, we noted that the stretch-out of the subthreshold drain current region was observed for only the  $\text{SiO}_2$  device, as shown in Fig. 3(b), suggesting that the light and bias combination stress can create some trap states at the interface between the IGZO and  $\text{SiO}_2$ . The reason that the  $\text{SiN}_x$  device did not exhibit such behavior is not clear at this time.

#### IV. CONCLUSION

An IGZO TFT with a PECVD-derived  $\text{SiO}_2$  gate dielectric exhibited greater bias stability and light-induced bias stability than the device with the  $\text{SiN}_x$  gate dielectric. This phenomenon

has been discussed based on a charge trapping mechanism. Therefore, the  $\text{SiO}_2$ -gated IGZO TFTs can be implemented as backplane electronics in the next-generation AM flat panel displays.

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