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# Effect of hydrogen on dynamic charge transport in amorphous oxide thin film transistors

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## Abstract

Hydrogen in zinc oxide based semiconductors functions as a donor or a defect de-activator depending on its concentration, greatly affecting the device characteristics of oxide thin-film transistors (TFTs). Thus, controlling the hydrogen concentration in oxide semiconductors is very important for achieving high mobility and minimizing device instability. In this study, we investigated the charge transport dynamics of the amorphous semiconductor InGaZnO at various hydrogen concentrations as a function of the deposition temperature of the gate insulator. To examine the nature of dynamic charge trapping, we employed short-pulse current–voltage and transient current–time measurements. Among various examined oxide devices, that with a high hydrogen concentration exhibits the best performance characteristics, such as high saturation mobility ( $10.9 \text{ cm}^2 \text{ v}^{-1} \text{ s}^{-1}$ ), low subthreshold slope ( $0.12 \text{ V/dec}$ ), and negligible hysteresis, which stem from low defect densities and negligible transient charge trapping. Our finding indicates that hydrogen atoms effectively passivate the defects in subgap states of the bulk semiconductor, minimizing the mobility degradation and threshold voltage instability. This study indicates that hydrogen plays a useful role in TFTs by improving the device performance and stability.

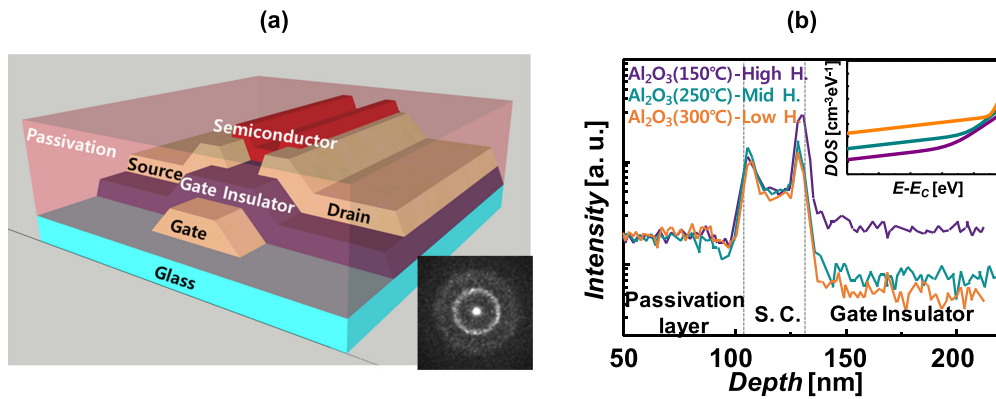
Keywords: oxide semiconductor, charge transport, InGaZnO, hydrogen, defects

(Some figures may appear in colour only in the online journal)

Hydrogen in zinc oxide based semiconductors, represented by InGaZnO (IGZO), ZnSnO (ZTO) and InHfZnO (IHZO), have emerged as promising candidates for use in switching and driving devices in active-matrix liquid crystal displays as well as wearable devices owing to their strong advantages, such as large-area device integration, low-cost fabrication process and perfect process compatibility with previous display technologies, which stems from their outstanding electrical, material, and optical properties [1–5]. The amorphous structure in bulk semiconductors has a certain number of inherent subgap densities of states, which determine both the electrical characteristics and reliability [6–8]. In addition, because of the low-temperature process used in this technology, interfacial defects between the gate insulator and semiconductor exist in

TFT devices, affecting the device mobility and stability [9, 10]. To obtain high mobility and high stability, various approaches to defect engineering have been introduced using various processes such as cation or anion optimization; oxygen vacancy manipulation; bi-, tri-, or multilayer adoption; and post photo/annealing processes [11–15]. However, precise control of defects in oxide semiconductor TFTs remains problematic owing to a significant number of subgap states in these semiconductors [16–20]. Thus, it is necessary to minimize defects in oxide TFTs as well as to understand their implications for device performance.

It is well known that hydrogen can be used to tune the electrical properties of oxide semiconductors [21–24]. It has a tendency to bond with oxygen in oxide semiconductors



**Figure 1.** (a) Schematic structure of a bottom gate *a*-IGZO TFT. (b) Data of the SIMS depth profile for hydrogen concentration in *a*-IGZO with increasing  $\text{Al}_2\text{O}_3$  deposition temperature. Inset shows density of states versus energy for  $\text{Al}_2\text{O}_3$  stacked *a*-IGZO TFT with hydrogen concentration.

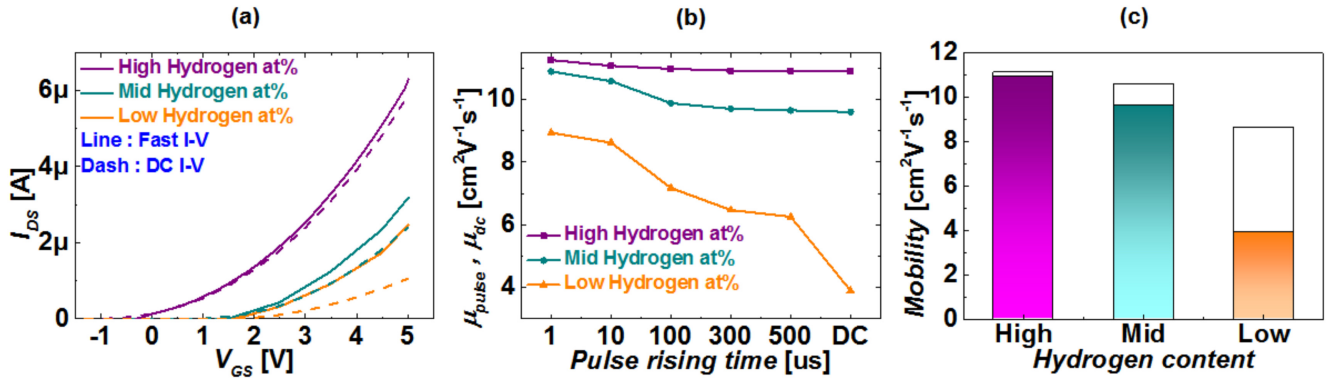
because of the very different electronegativities of these two elements [21–24]. Thus, hydrogen is always found in oxide semiconductors. It deactivates defects or generates shallow donors depending on the hydrogen concentration. When it forms bonds with oxygen and generates negatively charged hydroxide anions, hydrogen becomes a shallow donor in oxide semiconductors; consequently, a device becomes conductive and exhibits an insignificant on–off current ratio [21–24]. Very recently, however, various researchers have reported on beneficial effects of hydrogen, such as defect deactivation [25–27]. To date, however, the effect of hydrogen on defects in IGZO and on the TFT characteristics are not well examined or clearly understood. A critical reason is that the TFT characteristics are generally measured by conventional DC current–voltage ( $I$ – $V$ ) measurement, which provides an inherent charge-trapping environment due to the relatively long measurement time.

In this regard, microsecond-pulse  $I_D$ – $V_G$  measurement is quite useful for minimizing the effects of fast charge trapping, which causes a few unfavorable characteristics, such as threshold voltage instability and underestimation of the mobility of transistors [28–32]. In comparison to conventional DC measurement, which takes a few seconds, the microsecond short-pulse  $I_D$ – $V_G$  method takes a few microseconds to a few tens of microseconds, making it possible to minimize the effects of fast transient charging [28–32]. Consequently, the device characteristics can be measured under conditions with few or no charge traps, making it possible to evaluate the intrinsic device characteristics and to understand the dynamic charge transport on a time scale of microseconds [28–32].

In this paper, we examine the effect of the hydrogen concentration in oxide semiconductors on dynamic charge transport in oxide TFTs. To study the dynamic charge trapping and detrapping behavior as well as carrier transport mechanism, we employed microsecond short-pulse drain current–gate voltage ( $I_D$ – $V_G$ ) characterization and transient  $I_D$ –time measurements. We found that, in comparison to the mobility ( $\mu$ ) obtained by the DC  $I_D$ – $V_G$  method, the channel  $\mu$  values of TFTs obtained by fast pulse  $I_D$ – $V_G$  measurement increased by 1.6%–121.1% depending on the hydrogen

concentration in oxide semiconductors. In addition, our research results show that an oxide TFT with a high hydrogen concentration exhibits fewer defects and is less susceptible to ambient charge trapping, indicating that hydrogen plays a helpful role in defect passivation.

Bottom-gated TFTs of IGZO were fabricated using conventional thin film deposition, photolithography, and etching methods. A schematic diagram is presented in figure 1(a). A 40 nm thick IGZO thin film was prepared by radio-frequency sputtering using Ar and  $\text{O}_2$  gases. The composition ratio of In, Ga and Zn of the IGZO film were 1, 1, and 2.5, respectively. The structure of the IGZO thin film was examined by transmission electron microscopy. The results revealed tiny bright dots with ring patterns in a hazy background (figure 1(a), inset), indicating that the IGZO semiconductor material is formed of an amorphous medium. This amorphous IGZO semiconductor is referred to as *a*-IGZO hereafter. For all the electrodes, thin films of the transparent conductor indium tin oxide were deposited to a thickness of 100 nm by radio-frequency sputtering using Ar gas. Atomic layer deposition of the dielectric material  $\text{Al}_2\text{O}_3$  was performed using trimethylaluminum as the metal precursor and water as the reactant; this was followed by passivation and contact etching processes. The hydrogen concentration in *a*-IGZO was adjusted by increasing the deposition temperature of the  $\text{Al}_2\text{O}_3$  gate insulator from 150 °C to 250 °C and 300 °C. For hydrogen diffusion step, we performed annealing process at 300 °C in vacuum. This is quite effective manner of hydrogen diffusion into oxide semiconductor [21]. Secondary ion mass spectrometry was used to examine the hydrogen concentration. As seen in figure 1(b), the hydrogen concentration in *a*-IGZO decreases with increasing  $\text{Al}_2\text{O}_3$  deposition temperature. There is a tendency of hydrogen concentration in the front channel and front interface of  $\text{Al}_2\text{O}_3$ /IGZO as well as back interface of IGZO/ $\text{SiO}_2$  with  $\text{Al}_2\text{O}_3$  deposition temperature. The  $\text{Al}_2\text{O}_3$  stacked *a*-IGZO TFT processed at 150 °C exhibits a higher hydrogen concentration in the front channel and front/back interface than the TFTs processed at 250 °C and 300 °C. The relative concentrations of hydrogen for three samples in order from highest to lowest were 2, 1.3, and 1.0 respectively. This



**Figure 2.** (a) Fast  $I_D-V_G$  and DC  $I_D-V_G$  transfer curves of  $a$ -IGZO TFTs with various hydrogen concentrations. (b) Mobility values of  $a$ -IGZO TFTs with various hydrogen concentrations obtained by fast  $I_D-V_G$  with respect to pulse rising time and DC  $I_D-V_G$  measurement techniques. (c) Mobility values of  $a$ -IGZO TFTs with various hydrogen concentrations obtained by fast  $I_D-V_G$  (white) and DC  $I_D-V_G$  measurement (colored) techniques.

difference also affects the subgap states in  $a$ -IGZO which were extracted by a multifrequency capacitance–voltage ( $C-V$ ) characterization method [33, 34], as shown in the inset of figure 1(b). The  $a$ -IGZO TFT with a high hydrogen concentration shows lower subgap states than those with a low hydrogen concentration, indicating that hydrogen passivates the subgap states in  $a$ -IGZO [21]. Also Hall measurements data presents high hydrogen device presents high carrier concentration and high mobility in IGZO active layer, indicating that hydrogen plays in donor in oxide semiconductor.

To understand the effect of the hydrogen concentration in the  $a$ -IGZO semiconductor on the device characteristics of the  $a$ -IGZO TFT, we performed DC and microsecond fast-ramp  $I_D-V_G$  measurements using a semiconductor device parameter analyzer with a pulse generator and a high-resolution oscilloscope [28–32]. The DC  $I_D-V_G$  measurements were made at a scan speed of  $4 \text{ V s}^{-1}$ , and the fast-ramp  $I_D-V_G$  measurements were made at a voltage ramping speed of  $0.1\text{--}50 \mu\text{s V}^{-1}$ . Figure 2(a) shows the DC and microsecond fast-ramp  $I_D-V_G$  curves (voltage rise time of  $1 \mu\text{s}$ ) of  $a$ -IGZO TFT devices with various hydrogen concentrations. The  $a$ -IGZO TFT with a high hydrogen concentration exhibits a high  $\mu$  value ( $17.9 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ) and low subthreshold slope ( $s$ ) ( $0.11 \text{ V/dec.}$ ), whereas those with intermediate and low hydrogen concentrations exhibit values of  $\mu = 8.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $S = 0.17 \text{ V/dec}$  and  $\mu = 2.2 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $S = 0.21 \text{ V/dec}$  respectively. High mobility characteristics of the  $a$ -IGZO TFT with a high hydrogen concentration can be explained by the combined effect of low defect and high carrier concentration. For all the tested devices, the microsecond fast-ramp  $I_D-V_G$  curves show a higher  $I_D$  level and steeper slope than the DC  $I_D-V_G$  curves. As seen in figure 2(b), the mobility values of the  $a$ -IGZO TFT device with a high hydrogen concentration decrease gradually with increasing voltage rise time, whereas those of the TFTs with intermediate and low hydrogen concentrations decrease significantly with increasing the ramp time. Figure 2(c) shows the  $\mu$  values of  $a$ -IGZO TFTs with various hydrogen concentrations extracted from the DC and fast-ramp  $I_D-V_G$  curves. As expected, in agreement with figures 2(a) and (b), the  $\mu$  values extracted from the fast-ramp

$I_D-V_G$  curves are higher than those taken from the DC  $I_D-V_G$  data, indicating that the effect of fast charge trapping can be suppressed in the fast-ramp  $I_D-V_G$  curve. Furthermore, the  $a$ -IGZO TFT with a low hydrogen concentration shows a significantly higher mobility increase rate ( $\mu_{\text{fast-pulse}}/\mu_{\text{DC}}$ ) than that with a high hydrogen concentration, suggesting that the former is susceptible to fast charge trapping due to high subgap states in the semiconductor, as shown in the inset of figure 1(b). The  $\mu$  values of the  $a$ -IGZO TFT devices with high, intermediate, and low hydrogen concentrations increased by 1.58%, 10.25%, and 121.1%, respectively, after the fast ramp  $I_D-V_G$  measurement.

To further understand the effect of the hydrogen concentration in  $a$ -IGZO semiconductors on the charge trapping and detrapping characteristics of  $a$ -IGZO TFT devices, the hysteresis characteristics were measured using microsecond short-pulse  $I_D-V_G$  curves at a pulse width of 2 ms. The test setup is illustrated in figure 3(a). Figure 3(b) shows the microsecond single-pulse  $I_D-V_G$  curves (obtained at a pulse rise/fall times of  $1 \mu\text{s}$  and a pulse width of 4 ms) for three  $nc$ -IGZO TFTs. As shown on the left side of figure 3(b), we obtained two sets of  $I_D-V_G$  curves measured during the pulse rise/fall times. During the pulse fall time, the lower current level in the  $I_D-V_G$  data can be explained by fast charge trapping during the pulse duration [28–32]. During the pulse duration of 2 ms, the  $a$ -IGZO TFTs with high, intermediate, and low hydrogen concentrations exhibit  $\Delta V_{\text{th}}$  values of 0.41, 1.58, and 3.34 V, indicating that  $1.81 \times 10^{11}$ ,  $6.94 \times 10^{11}$ , and  $1.59 \times 10^{12} \text{ cm}^{-2}$  carriers, respectively, were trapped in the subgap states in the semiconductor. The right side of figure 3(b) shows the transient  $I_D$ -time characteristics of  $a$ -IGZO TFT devices with various hydrogen concentrations, which illustrate fast transient charge trapping during the pulse duration. As the gate bias voltage in the gate pulse reached the plateau region, the  $I_D$  levels started to decrease. The reduction in the  $I_D$  levels was most noteworthy in the  $a$ -IGZO TFT with the low hydrogen concentration, suggesting that this device has a significant number of subgap defects in the interface and bulk oxide semiconductor as compared with the  $a$ -IGZO TFT with a high hydrogen concentration.

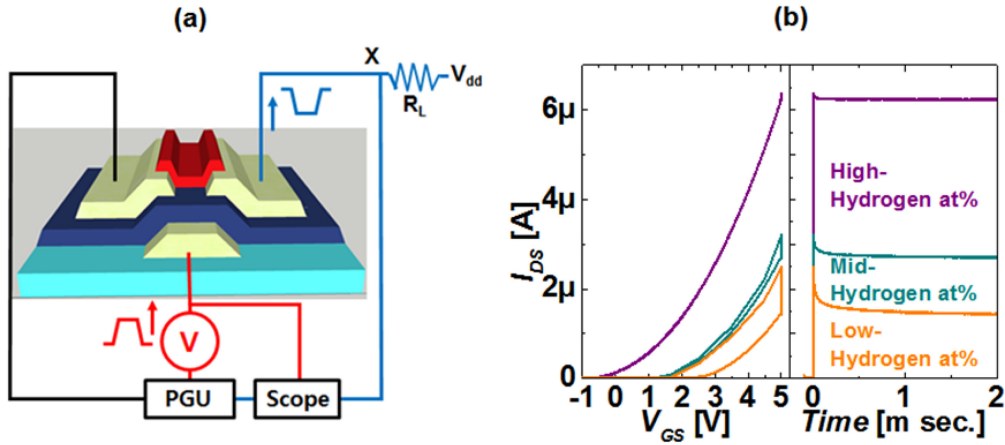


Figure 3. (a) Schematics of pulse  $I_D$ - $V_G$  measurement system setup. (b) Pulse  $I_D$ - $V_G$  curves and corresponding transient  $I_D$ -time data of  $a$ -IGZO TFTs with various hydrogen concentrations.

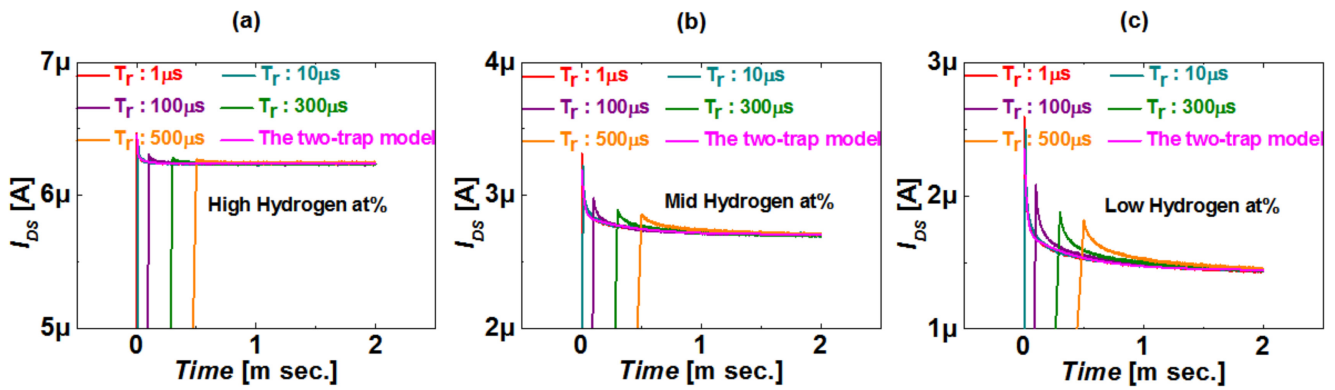


Figure 4. Transient  $I_D$ -time characteristics of  $a$ -IGZO TFTs with various hydrogen concentrations as a function of pulse rising time.

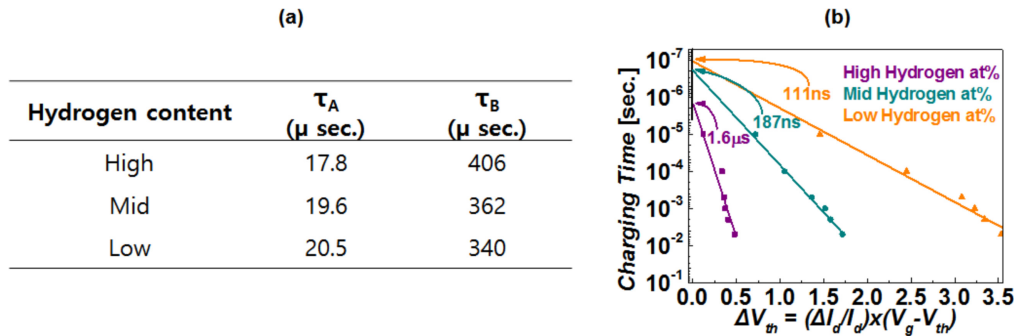
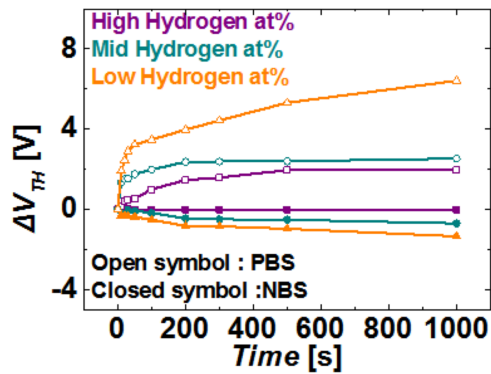


Figure 5. (a) Time constants of  $a$ -IGZO TFTs with various hydrogen concentrations extracted from two trap fitting model. (b)  $\Delta V_{th}$ -charging time of  $a$ -IGZO TFTs with various hydrogen concentrations.

To assess the dynamic charge trapping properties of the  $a$ -IGZO TFTs, we measured the transient  $I_D$ -time curves for rise times of 1–500  $\mu$ s, as seen in figures 4(a)–(c). During the pulse duration,  $I_D$  decayed rapidly because of fast transient charge trapping, followed by gradual saturation within a few milliseconds. Then,  $I_D$  approached a certain value with a typical charge trapping curve. The measured transient  $I_D$ -time data were fitted with two-trap models using the equation  $I = AI_0 \exp(-t/\tau_A) + BI_0 \exp(-t/\tau_B)$  with two time constants [35]. The best-fitted one and the measured data are presented in figures 4(a)–(c) for the  $a$ -IGZO TFTs with high,

intermediate, and low hydrogen concentrations, respectively. The two-trap model for the  $a$ -IGZO TFTs with various hydrogen concentrations described well all the charge-trapping regions at the initial, intermediate, and long time scales. Note that the transient  $I_D$  variation with time obviously follows multiple trap-to-trap charging processes. Two different time constants for the  $a$ -IGZO TFTs are presented in figure 5(a). The shorter trap time constants ( $t_A$ ) for three devices are similar while the longer trapping time constants ( $t_B$ ) are significantly influenced by hydrogen concentration in oxide TFTs.



**Figure 6.** Positive bias stress (open symbol) and negative bias stress (closed symbol) characteristics of *a*-IGZO TFTs with various hydrogen concentrations.

To quantitatively analyze the trap time constant, we estimated the threshold voltage shift ( $\Delta V_{th}$ ) from the measured transient  $I_D$ -time data using the equation  $\Delta V_{th} = \Delta I_D(V_G - V_{th})/I_D$ , where  $\Delta I_D$  is the difference between the  $I_D$  values at the initial and final points of the pulse duration,  $V_G$  is the pulse amplitude,  $V_{th}$  is the threshold voltage and  $I_D$  is the maximum drain current at the pulse onset [28–31]. The charging time versus  $\Delta V_{th}$  is plotted for the *a*-IGZO devices with various hydrogen concentrations in figure 5(b). The critical trapping time constant ( $\tau_C$ , the time at which charge trapping begins) can be estimated from the curve of the charging time versus  $\Delta V_{th}$ . The *a*-IGZO TFT with a high hydrogen concentration exhibits a long  $\tau_C$  of 1.6  $\mu$ s, whereas the  $\tau_C$  values of the *a*-IGZO TFTs with intermediate and low hydrogen concentrations are 187 and 111 ns, respectively.

To probe whether charge transport occurs primarily in the interface or the bulk, we measured the low-frequency noise (LFN). According to the normalized power spectral density  $S_{ID}$  versus  $|V_{GS} - V_{th}|^{-x}$  for  $x \sim 1.03$  (high hydrogen concentration),  $x \sim 1.12$  (intermediate hydrogen concentration), and  $x \sim 1.399$  (low hydrogen concentration), which are close to 1, carrier transport occurs mainly via the bulk channel [36–39]. Thus, hydrogen, which effectively deactivates bulk subgap states in semiconductors, improves the device performance and reliability of *a*-IGZO TFTs [36–39]. Therefore, as we determined in positive and negative bias stability (PBS and NBS) tests, (PBS,  $V_G = 20$  V,  $V_D = 0$  V, NBS:  $V_G = -20$  V,  $V_D = 0$  V), the *a*-IGZO TFT with a high hydrogen concentration exhibits a lower  $\Delta V_{th}$  (figure 6), which is related to the subgap states in the front channel.

In this paper, we described the dynamic charge transport in *a*-IGZO TFTs with various hydrogen concentrations. To this end, we employed fast-ramp  $I_D$ - $V_G$ , short-pulse  $I_D$ - $V_G$ , and transient  $I_D$ -time measurements. Further, to probe the subgap states in *a*-IGZO, we applied multifrequency  $C$ - $V$  and LFN methods. Among oxide TFTs with various hydrogen concentrations, the device with a high hydrogen concentration exhibited the best performance (high mobility and low subthreshold slope) and negligible charge trapping

characteristics, which resulted from the low subgap states in *a*-IGZO. This indicates that the incorporation of hydrogen into *a*-IGZO deactivates the subgap defects in the bulk semiconductor, minimizing the mobility degradation and threshold voltage instability. This study demonstrates that hydrogen plays a useful role in improving the device performance and reliability of oxide TFTs. Therefore, controlling the hydrogen concentration is crucial for successful application of oxide TFTs.

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